

RESEARCH ARTICLE

Control Variates method to estimate stochastic buckling loads

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Funding Information

This research was supported by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) in the framework of project 511267658

Abstract

Buckling is the most significant failure mode for thin-walled structures. In particular, geometric imperfections have a major influence on the buckling behavior. These spatial correlated imperfections are inherently random and can be modeled using random fields. Therefore, computationally expensive probabilistic buckling analyses have to be performed. For some structures, a linear pre-buckling behavior can be observed. In this case, the stability point can be calculated with a linear buckling analysis, which is widely used in engineering. However, the results of linear buckling analyses strongly differ from the correct buckling load in case of a non-linear pre-buckling behavior. Then, a non-linear buckling analysis is required, which is computationally expensive for probabilistic safety assessments based on Monte Carlo simulations.

This paper aims to estimate the second-order statistics of buckling loads for thin-walled structures exhibiting strongly non-linear pre-buckling behavior. The estimation leverages existing correlations between the outcomes of linear and non-linear buckling analyses. The proposed approach utilizes the framework of Control Variates, wherein the more expensive analysis (non-linear buckling analysis) is run a few times only, while the cheaper linear buckling analysis is run a considerable number of times. The proposed method is demonstrated on a variety of structures, including a folded plate with multiple types of stability points, a composite shell panel and a cylinder with random geometric imperfections. In these numerical examples, stochastic buckling analysis using Control Variates is approximately 5 to 20 times faster than classical Monte Carlo simulation.

KEYWORDS

Buckling analysis, Monte Carlo simulation, Control Variates, Second-order statistics, Random imperfections Aleatory uncertainty

1 | INTRODUCTION

One of the dominant failure modes of thin-walled structures is buckling. Geometric and material imperfections, such as deviations in shape and thickness, residual stresses, variations in boundary conditions and material properties have a substantial influence on the buckling behavior. Even small variations in geometric imperfections significantly influence the load-bearing capacity. The exact shape of imperfections is frequently unknown or in other words uncertain. Deterministic and semi-probabilistic design concepts are based on very conservative design factors, also known as knockdown factors (KDFs). The development of reliable and more economical KDFs remains a focus of numerous ongoing research projects, as illustrated, e.g., in [38].

In a probabilistic approach, the aleatory uncertainties of spatially varying imperfections are modeled as random fields, see, e.g., in [31, 7, 26]. In this paper, random imperfections are applied to the finite element (FE) model, and Monte Carlo simulations (MCS) are conducted to determine the second-order statistics of random buckling loads. However, to quantify the uncertain shape of the geometrical imperfections only a few measurements are available. The definition of deterministic random field parameters under these limited data would imply a precise probabilistic knowledge on the stochastic distribution and its spatial correlation. Therefore, in [17, 14, 18], the concept of polymorphic (mixed/hybrid) uncertainty models [23], also known as imprecise probabilities [6, 12], is introduced to consider epistemic uncertainties of shell imperfections by means of

intervals or fuzzy numbers. An extension to quantify the uncertainties in boundary conditions, material properties, and thickness imperfections is provided in [16]. This paper initially explores the feasibility of using the Control Variates method for stochastic buckling analyses, with an emphasis on the spatial correlation of random imperfections. A potential extension of this work could involve applying the method to address polymorphic uncertainty, such as random fields with interval or fuzzy correlation lengths.

A numerical treatment of aleatory and epistemic uncertainties requires a multi-loop algorithm consisting of the fundamental solution (e.g. FE buckling analysis), the MCS loop and the fuzzy or interval analysis. This can be highly computationally expensive and the effort increases rapidly for buckling design optimization, as illustrated in [15]. Additionally, when geometric imperfections are applied to an FE model, the geometrical deviation at each node is a single input variable. This leads to a high-dimensional input space of uncertain variables. Therefore, appropriate surrogate models can be used to replace the time consuming FE analysis. For instance, an approach for efficiently analyzing the imperfection sensitivity using reduced order models is presented in [4, 22]. Furthermore, neural network surrogate models to approximate FE buckling analysis are shown in [24, 32]. Based on a first surrogate model for the fundamental solution, a further surrogate model can be constructed to replace the MCS, see for example [21]. An approach for bounding imprecise failure probabilities of linear structural systems is introduced in [19]. However, this approach is not applicable to non-linear buckling problems, as it relies on the assumption of a linear mapping between input and output quantities within the used operator norm theory.

This paper aims to estimate the second-order statistics (mean and standard deviation) of the buckling loads of thin-walled shell structures, with imperfections characterized by probabilistic models. The proposed approach is based on the concept of Control Variates [3], which has been studied in various engineering applications, such as [1, 8, 30]. The main contribution of this work is to introduce Control Variates for probabilistic buckling analysis, where geometric imperfections are modeled as random fields. The idea is to leverage the existing correlations between the solutions of linear and non-linear buckling analyses.

In case of non-linear pre-buckling behavior, a non-linear buckling analysis provides an accurate prediction of the buckling load, and a linear buckling analysis leads to imprecise results. It should be noted, that for some structures with a linear pre-buckling behavior, the linear buckling analysis can still yield an incorrect prediction due to buckling mode interaction. However, in case of multiple buckling loads nearly at the same level, e.g., for an axially loaded cylinder, the "exact" buckling load can be calculated and the shape of the associated buckling mode is not relevant. Furthermore, the non-linear buckling analysis requires a geometrically non-linear path-following analysis involving an iterative procedure. Therefore, the linear buckling approach is typically less computationally expensive than the non-linear buckling analysis. Second-order statistics are determined by sampling, involving a limited number of non-linear analyses and a relatively large number of linear buckling analyses. Thus, the presented Control Variates approach reduces overall computational costs, as the more expensive non-linear buckling analysis is performed only a few times, while the less costly linear buckling analysis is executed more frequently. Furthermore, by exploiting the correlation between linear and non-linear analyses, it is still possible to estimate the statistics of the accurate buckling load, even when the linear buckling analysis is applied to cases with strong non-linear pre-buckling behavior.

In this paper, the Control Variates approach is demonstrated for stochastic buckling analysis of a composite cylindrical shell panel. Random geometric imperfections are modeled as random fields using the Karhunen-Loève Expansion (KLE). The shape of the random imperfection can be controlled by the correlation length. Thus, second-order statistics of the buckling loads are analyzed for various correlation lengths using Monte Carlo simulations (MCS). To illustrate the applicability of Control Variates, the correlation between linear and non-linear buckling analysis is investigated. The effectiveness of Control Variates in estimating the second-order statistics of buckling loads is demonstrated on various thin-walled structures, including a folded plate with different types of stability points, a composite shell panel and a cylinder subjected to random geometric imperfections with varying correlation lengths.

The paper's innovative contributions and key features can be summarized as follows:

- Introduction of Control Variates with splitting technique for stochastic buckling analyses
- Efficient estimation of the second-order statistics of buckling loads
- Leveraging correlations between linear and non-linear buckling analyses
- Random field modeling with Control Variates
- Study to the effectiveness of Control Variates for various correlated random geometric imperfections
- Stochastic buckling analysis using Control Variates is approximately 5 to 20 times faster than classical Monte Carlo simulation

In Sections 2 and 3, the fundamentals of numerical buckling analysis and the estimation of second-order statistics using Monte Carlo simulations are presented. Following this, the Control Variates approach for buckling is introduced in Section 4. Section 5

70 illustrates the concept on three different examples: a folded plate, a composite cylindrical shell panel and a composite cylinder.
71 Finally, Section 6 provides of conclusion and outlines potential directions for future research.

72 2 | BASICS OF LINEAR AND NON-LINEAR BUCKLING ANALYSIS

Various strategies are available for identifying a stability point, as discussed, e.g., in [43, 40]. For a specific load level $\lambda \mathbf{P}_0$ with a load factor λ and a basic external load \mathbf{P}_0 , a non-linear eigenvalue problem can be constructed

$$[\mathbf{K}_{\text{lin}} + \Lambda \mathbf{K}_{\text{nonlin}}] \boldsymbol{\varphi} = \mathbf{0}, \quad (1)$$

where it is assumed that

$$\mathbf{P}_{\text{cr}} \sim \Lambda \mathbf{K}_{\text{nonlin}}. \quad (2)$$

In Eq. (1), Λ can be interpreted as a load increasing factor and $\boldsymbol{\varphi}$ is the associated eigenvector. The tangent stiffness matrix \mathbf{K}_T is divided in linear \mathbf{K}_{lin} and non-linear parts $\mathbf{K}_{\text{nonlin}}$. Generally, the tangent stiffness matrix depends on the displacement \mathbf{u} and stress state $\boldsymbol{\sigma}(\mathbf{u})$, respectively. If the variational formulation allows to separate $\mathbf{K}_{\text{nonlin}}$, the initial displacement matrix \mathbf{K}_U and the geometrical matrix \mathbf{K}_G can be introduced

$$\mathbf{K}_T = \mathbf{K}_{\text{lin}} + \mathbf{K}_{\text{nonlin}} = \mathbf{K}_{\text{lin}} + \mathbf{K}_U(\mathbf{u}) + \mathbf{K}_G(\boldsymbol{\sigma}(\mathbf{u})). \quad (3)$$

A solution for $\Lambda = 1$ in Eq. (1) yields the classical form of an eigenvalue problem for the tangent stiffness matrix

$$(\mathbf{K}_{\text{lin}} + \Lambda \mathbf{K}_{\text{nonlin}}) \boldsymbol{\varphi} = \mathbf{0} \Leftrightarrow \mathbf{K}_T \boldsymbol{\varphi} = \mathbf{0} \Leftrightarrow (\mathbf{K}_T - \omega \mathbf{1}) \boldsymbol{\varphi} = \mathbf{0}, \quad (4)$$

wherein a stability point is indicated for $\omega = 0$. In the non-linear case, the eigenvalue Λ is an indicator for the type of stability

$$\begin{aligned} \Lambda > 1 &\rightarrow \text{stable} \\ \Lambda = 1 &\rightarrow \text{indifferent (stability point)} \\ \Lambda < 1 &\rightarrow \text{unstable.} \end{aligned} \quad (5)$$

Furthermore, the non-linear critical load vector can be computed by

$$\mathbf{P}_{\text{cr}} = \mathbf{P}_{\text{cr,nonlin}} = \Lambda(\lambda \mathbf{P}_0) \quad \text{with} \quad \Lambda = 1. \quad (6)$$

In addition, the type of stability point can be determined by the following criterion, see, e.g., [33, 44],

$$\boldsymbol{\varphi}^T \mathbf{P}_{\text{cr}} \begin{cases} = 0 & \text{bifurcation point} \\ \neq 0 & \text{limit point.} \end{cases} \quad (7)$$

73 A schematic non-linear load-displacement curve $\lambda \mathbf{P}_0 - u$ with corresponding eigenvalue-displacement curves $\Lambda(\lambda \mathbf{P}_0) - u$ and
74 $\omega - u$ are depicted in Fig. 1, see black, red and blue curve, respectively.

For some structures, a linear pre-buckling behavior can be observed. In such cases, only a single linear calculation step and the solution of an eigenvalue problem are required. This motivates the use of the linear buckling analysis as a special case of the non-linear analysis. It starts from the displacement state $\mathbf{u} = \mathbf{0}$, where the linear solution

$$\mathbf{K}_T(\mathbf{0}) \mathbf{u}_0 = \mathbf{P}_0 \quad \Leftrightarrow \quad \mathbf{u}_0 = \mathbf{K}_T^{-1}(\mathbf{0}) \mathbf{P}_0 \quad (8)$$

is computed for an external basic load of \mathbf{P}_0 ($\lambda_0 = 1$) with $\mathbf{K}_T(\mathbf{0}) = \mathbf{K}_{\text{lin}}$. Thus, the linear buckling analysis is defined as

$$[\mathbf{K}_{\text{lin}} + \Lambda_0 \mathbf{K}_{\text{nonlin}}(\mathbf{u}_0)] \boldsymbol{\varphi}_0 = \mathbf{0}. \quad (9)$$

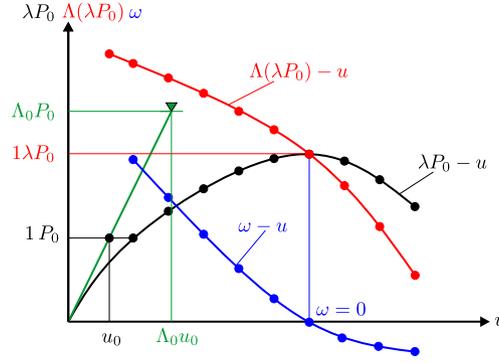


FIGURE 1 Schematic representation of linear and non-linear buckling analysis

The associated critical load and displacement vectors are

$$\tilde{\mathbf{P}}_{\text{cr}} = \mathbf{P}_{\text{cr,lin}} = \Lambda_0 \mathbf{P}_0, \quad (10)$$

$$\tilde{\mathbf{u}}_{\text{cr}} = \mathbf{u}_{\text{cr,lin}} = \Lambda_0 \mathbf{u}_0. \quad (11)$$

75

However, as shown in Fig. 1, the results of a linear buckling analysis (maximum value of the green curve) can significantly differ from the non-linear buckling load (maximum value of the black curve) in case of non-linear pre-buckling behavior. This requires a comprehensive geometrically non-linear path-following analysis using an iterative procedure such as the Newton-Raphson scheme. In this paper, the non-linear buckling analysis is performed by a path-following analysis, where the signs of the diagonal elements of the tangent stiffness matrix are observed. A change in the sign of the diagonal elements of \mathbf{K}_T indicates a change in the equilibrium state, with

$$\begin{aligned} \forall D_{ii}, D_{ii} > 0 &\rightarrow \text{stable} \\ \exists D_{ii}, D_{ii} = 0 &\rightarrow \text{indifferent, (stability point).} \\ \exists D_{ii}, D_{ii} < 0 &\rightarrow \text{unstable.} \end{aligned} \quad (12)$$

76 If at least one of the diagonal elements D_{ii} becomes negative, the calculation is terminated and the load state is saved. At this
77 equilibrium state, the critical load vector \mathbf{P}_{cr} and the initial post-buckling mode φ_{cr} can be computed with the non-linear
78 eigenvalue problem given by Eq. (1). Based on the special loading conditions in the examples, the buckling loads P_{cr} and \tilde{P}_{cr} can
79 be computed from the associated load vectors \mathbf{P}_{cr} and $\tilde{\mathbf{P}}_{\text{cr}}$.

80 3 | SECOND-ORDER STATISTICS OF BUCKLING LOAD

81 3.1 | Uncertainty in buckling load

82 The behavior of structures prone to buckling may be affected by several factors, which cannot be quantified deterministically. For
83 example, there may be imperfections in the shells with respect to their nominal dimensions due to uncertainty associated with the
84 manufacturing processes. Also parameters such as Young's modulus or Poisson's ratio may be affected by uncertainties due to
85 variability associated with batches of raw material. Therefore, these sources of uncertainty must be explicitly accounted for such
86 that their effect can be properly quantified. In the following, it is assumed that material and geometrical properties of a structure
87 are collected in the vector ξ . The uncertainty associated with ξ is characterized by a random variable vector Ξ with probability
88 density function $p_{\Xi}(\xi)$ [11]. As the buckling load of a structure P_{cr} depends on the properties of the structure ξ , which are
89 assumed to be uncertain, it is clear that the buckling load becomes uncertain as well. In other words, the buckling load P_{cr} is
90 a random variable, with its own probability density function. Calculating the probability density function associated with the
91 buckling load may be challenging for problems of engineering interest, particularly in the tails of the distribution. However, the
92 calculation of second-order statistics (that is, mean and variance) may be more tractable while still providing valuable insights
93 on the uncertainty associated with the buckling load.

3.2 | Estimation of second-order statistics of buckling loads by means of Monte Carlo simulation

Recall that the non-linear buckling load P_{cr} depends on the properties of the structure ξ . For cases of practical interest, there exists no closed-form relation between the model properties and buckling load. Rather, they must be derived from an FE analysis. As such, an analysis is usually only available in the form of a black-box. That is, for a given input ξ , one obtains the output P_{cr} . In this context, Monte Carlo simulation appears as a natural alternative for computing the sought second-order statistics, see e.g. [20]. In a nutshell, a Monte Carlo simulation consists of generating n independent samples $\xi^{(j)}$, $j = 1, \dots, n$ of the properties of the structure distributed according to $p_{\Xi}(\xi)$. Then, the buckling loads for each of the samples are calculated, yielding $P_{cr}(\xi^{(j)})$, $j = 1, \dots, n$. Under the assumption that the samples of the properties $\xi^{(j)}$, $j = 1, \dots, n$ are grouped in matrix Ξ_n , estimates of the mean value $\widehat{\mu}'_1(P_{cr}, \Xi_n)$ and the variance $\widehat{\mu}_2(P_{cr}, \Xi_n)$ are obtained by the following expressions

$$\widehat{\mu}'_1(P_{cr}, \Xi_n) = \frac{1}{n} \sum_{j=1}^n P_{cr}(\xi^{(j)}), \quad (13)$$

$$\widehat{\mu}_2(P_{cr}, \Xi_n) = \frac{1}{n-1} \sum_{j=1}^n (P_{cr}(\xi^{(j)}) - \widehat{\mu}'_1(P_{cr}, \Xi_n))^2, \quad (14)$$

As these estimates are produced by Monte Carlo simulations, they are affected by inherent randomness associated with the sampling process. A means to quantify the quality of these estimates is by calculating their respective variances $\widehat{\sigma}^2[\cdot]$, see e.g. [2],

$$\widehat{\sigma}^2[\widehat{\mu}'_1(P_{cr}, \Xi_n)] = \frac{\widehat{\mu}_2(P_{cr}, \Xi_n)}{n}, \quad (15)$$

$$\widehat{\sigma}^2[\widehat{\mu}_2(P_{cr}, \Xi_n)] = \frac{\widehat{\mu}_4(P_{cr}, \Xi_n)}{n} - \frac{(n-3)\widehat{\mu}_2^2(P_{cr}, \Xi_n)}{(n-1)n}. \quad (16)$$

In the last equation, $\widehat{\mu}_4(P_{cr}, \Xi_n)$ denotes the estimator of the fourth-order central moment of the non-linear buckling load. It is estimated using Eq. (A3) in Appendix A.

Equations (15) and (16) indicate that the variances of the estimators for the second-order statistics depend on the number of samples n . It is desirable that these variances are as small as possible, as this increases the confidence that the quantities being estimated are good approximations of the exact second-order statistics. In other words, it is desirable to obtain estimates of these statistics with sufficient *precision*. However, small variances (or equivalently, high precision) may entail a large number of simulations n , which can be quite costly from a numerical viewpoint, as it implies performing n non-linear buckling analyses. Therefore, in practical applications, it is expected that n is actually small and therefore, the estimators will possess large variability. Such a concept is illustrated schematically in Fig. 2, where the probability density associated with an estimator $\widehat{\mu}$ (which represents either mean or variance, see red curve) is relatively flat, reflecting high uncertainty on the true value of the sought statistic.

4 | CONTROL VARIATES: A TOOL FOR AGGREGATING ESTIMATES OF NON-LINEAR AND LINEAR BUCKLING ANALYSES

4.1 | Control Variates

The objective is to estimate the second-order statistics of the non-linear buckling load of a structural system. Estimating these statistics with sufficient precision by means of Monte Carlo simulation may demand a considerable number of non-linear buckling analyses. However, from the discussion in Section 2, it is known that calculating linear buckling loads is numerically less demanding than its non-linear counterpart, although it leads to different values of the sought buckling load. A natural question in this scenario is: is it possible to leverage on linear analysis to estimate statistics of the non-linear buckling load? The answer to this question is affirmative, as this can be carried out within the framework of Control Variates [20, 29]. Indeed, let μ denote a statistic, which can represent either the mean μ'_1 or the variance μ_2 . Then, the estimator of μ considering Control Variates (CV) is the following:

$$\hat{\mu}^{(CV)} = \hat{\mu}(P_{cr}, \Xi_n) - \gamma \hat{\mu}(\tilde{P}_{cr}, \Xi_n) + \gamma \hat{\mu}(\tilde{P}_{cr}, \Xi_m). \quad (17)$$

In the above equation, Ξ_n and Ξ_m denote two sets with n and m samples, respectively; γ is the so-called *control parameter*, which is actually a real number whose calculation is discussed later on; and $\hat{\mu}(x, \Xi_y)$ denotes estimation of the statistics μ considering the response x (where x could represent either the non-linear or linear buckling load) and Ξ_y denotes a sample set (where y could represent either n or m). It is assumed that $m > n$. The idea behind the Control Variates estimator in Eq. (17) is the following.

- The last term $\gamma \hat{\mu}(\tilde{P}_{cr}, \Xi_m)$ denotes the sought statistic calculated considering the linear buckling load employing a large number of samples m , which is moreover amplified by γ . This last term by itself does not lead to the sought statistic because it involves the linear buckling load (amplified by γ) instead of its non-linear counterpart. However, as m is large, this estimator should possess a relatively low variance. This is illustrated schematically in Fig. 2, where the probability density associated with this estimator (shown with the green line) is quite peaked.
- The difference $\hat{\mu}(P_{cr}, \Xi_n) - \gamma \hat{\mu}(\tilde{P}_{cr}, \Xi_n)$ in Eq. (17) can be interpreted as a correction term, as it subtracts the value of the statistic associated with the linear buckling load amplified by γ and adds the statistic calculated with the non-linear buckling load. Usually, the variance of this difference should be relatively small, even if n itself is small. The reason is that while the linear and non-linear buckling loads produce different results, it is expected nevertheless that there is a high degree of correlation between them. The probability density associated with this difference is represented schematically with a violet line in Fig. 2.

The summation of the two terms described above leads to the Control Variates estimator of Eq. (17). In essence, such an estimator allows to *aggregate* the results stemming out of linear and non-linear buckling analyses but still leads to conclusions about the non-linear buckling load. In fact, the effect of the linear buckling load cancels out from Eq. (17), as the subtraction between the second and third terms of the right-hand side of that equation is equal to zero. But in that process, the presence of those second and third terms helps in decreasing the variance of the statistics of the buckling load. Indeed, this estimator usually possesses a relatively small variance, as illustrated schematically in Fig. 2 with the probability density function in blue color.

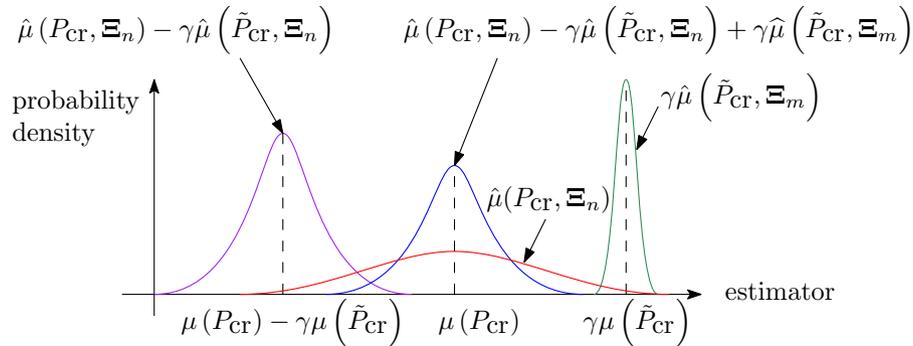


FIGURE 2 Schematic representation of Control Variates estimator.

The advantage of the Control Variates estimator is that in the calculation process, correlations between linear and non-linear buckling loads are exploited. This implies that the information retrieved from linear buckling analyses may improve the conclusions that are drawn respecting non-linear buckling analyses. In this sense, the aforementioned control parameter γ plays a key role. To understand this concept, it is essential to study the variance of the Control Variates estimator. It can be shown that the variance $\hat{\sigma}^2$ associated with the Control Variates estimator is equal to [20]:

$$\hat{\sigma}^2 [\hat{\mu}^{(CV)}] = \hat{\sigma}^2 [\hat{\mu}(P_{cr}, \Xi_n)] - 2\gamma \hat{\sigma} [\hat{\mu}(P_{cr}, \Xi_n), \hat{\mu}(\tilde{P}_{cr}, \Xi_n)] + \gamma^2 \hat{\sigma}^2 [\hat{\mu}(\tilde{P}_{cr}, \Xi_n)] + \gamma^2 \hat{\sigma}^2 [\hat{\mu}(\tilde{P}_{cr}, \Xi_m)], \quad (18)$$

153 where $\widehat{\delta}[\cdot, \cdot]$ denotes the covariance estimator between the arguments. It is noted that the variance of the Control Variates
 154 estimator as shown in Eq. (18) is a quadratic function with respect to the control parameter γ . Therefore, γ can be selected such
 155 that this variance is minimized, which implies forcing the derivative of Eq. (18) with respect to γ to be equal to zero. In such
 156 way, one determines the optimal control parameter γ^* , which is equal to:

$$\gamma^* = \frac{\widehat{\delta}[\widehat{\mu}(P_{cr}, \Xi_n), \widehat{\mu}(\check{P}_{cr}, \Xi_n)]}{\widehat{\sigma}^2[\widehat{\mu}(\check{P}_{cr}, \Xi_n)] + \widehat{\sigma}^2[\widehat{\mu}(P_{cr}, \Xi_m)]}. \quad (19)$$

157 In summary, the application of Eqs. (17), (18) and (19) allows estimating the sought statistic (either mean or variance) by means
 158 of Control Variates. These equations are applied in the following order. First, the buckling loads considering linear (\check{P}_{cr}) and
 159 non-linear analysis (P_{cr}) are calculated for each sample contained in the set Ξ_n . In addition, the buckling loads considering linear
 160 analysis only are evaluated for each sample contained in the set Ξ_m . Considering all of these samples, Eq. (19) is evaluated to
 161 obtain the optimal control parameter γ^* . This optimal control parameter is then used together with the samples of the buckling
 162 loads associated with the sets Ξ_n and Ξ_m to evaluate the sought statistic through Eq. (17) as well as the variance of this estimator
 163 by means of Eq. (18).

164 As noted from the above description, the application of Control Variates is completely non-intrusive. That is, it is not necessary
 165 to access to system's matrices. It just suffices to conduct linear and non-linear buckling analyses for different sets of samples.
 166 Then, the Control Variates estimate in Eq. (17) merges the information contained in these samples to produce an estimate of
 167 the sought statistic involving non-linear buckling analysis. Indeed, the role that the samples of the linear buckling load in the
 168 estimator is simply exploiting correlations to reduce the variance of the estimator of the sought statistic. In that sense, the
 169 optimal control parameter γ^* in Eq. (19) plays a pivotal role. To understand its role better, consider the case where there is a
 170 high covariance between buckling loads calculated using linear and non-linear buckling analysis. In such situation, performing
 171 linear buckling analysis is almost *as good as* performing non-linear buckling analysis, meaning that both types of analyses are
 172 capable of uncovering the effects of uncertainty in the buckling load. In such situation of high covariance, the control parameter
 173 γ^* will naturally approach to 1 and thus, the Control Variates estimator provides more importance to the information carried
 174 by the samples of the linear buckling load, which translates into an estimator of the sought statistic with reduced variance. In
 175 the (unlikely) event that the covariance $\widehat{\delta}$ between the linear and non-linear buckling load is zero, the control parameter itself
 176 becomes zero and the estimator of Eq. (17) reduces to its plain Monte Carlo counterpart. In such a case, as there is no covariance
 177 $\widehat{\delta}$, the Control Variates estimate cannot leverage on the information carried by the samples of the linear buckling load and
 178 therefore, the sought estimate is produced based on the samples of the non-linear buckling analysis only.

179 To summarize the concepts described above, the application of Control Variates for estimating buckling loads can be visualized
 180 as follows. First, a relatively small number of samples of the buckling load are generated considering non-linear buckling
 181 analysis. As the number of samples drawn is small, the estimators drawn out of those samples (such as mean and variance) will
 182 be most likely highly uncertain (implying that their variances are high). But then, if one performs additional sampling resorting
 183 to linear buckling analyses, it is possible to exploit covariance between linear and non-linear buckling analysis to reduce the
 184 variance of the estimators associated with the non-linear buckling load.

185 4.2 | Control Variates with Splitting

The previous section has presented the application of the Control Variates framework. For its practical implementation, note
 that the same samples of the linear and non-linear buckling load are used for both evaluating the optimal control parameter
 (see Eq. (19)) and the sought statistic (see Eq. (17)). However, such a strategy induces bias in the estimator of the statistic,
 as documented, e.g., in [29]. The effect of bias can be particularly notorious in case that the sample set Ξ_n possesses a small
 number n of samples, which is expected to be precisely the case in practical applications, because performing n non-linear
 buckling analyses is numerically demanding. A remedy to eliminate bias is to resort to a *Splitting* approach, as proposed in [3].
 This is a quite convenient scheme, as it does not demand performing additional buckling analyses (neither linear nor non-linear).
 The splitting approach consists of dividing the set of available samples into subsets. Then, the associated estimators (e.g. mean,
 variance, optimal control parameter) are estimated for each of these subsets. Finally, the estimators for the subsets are aggregated
 in such a way that bias is effectively eliminated. As discussed in [3], the number of subsets to be considered should be equal or
 larger than 3. However, a large number of subsets may increase the variance. Therefore, in this paper, the minimum number of 3

subsets is considered to implement the splitting approach.

How does the splitting approach work in practice? First, each of the sample sets Ξ_n and Ξ_m is partitioned into three subsets $\Xi_{m^*,k}$ and $\Xi_{n^*,k}$, where $k = 1, 2, 3$ and $n^* = n/3$ and $m^* = m/3$. Here, it is implicitly assumed that n and m are selected such that they are multiples of 3. For each subset k , the subset controller $\tau(k)$ is defined [3], as shown in Table 1. Once that the subsets

TABLE 1 Subset controllers $\tau(k)$ for each subset k .

Subset k	Subset controller $\tau(k)$
1	2
2	3
3	1

have been defined, the expressions for calculating the sought statistic, its variance and the optimal control parameter by means of Control Variates with Splitting (CV+S) are:

$$\widehat{\mu}^{(CV+S)} = \frac{1}{3} \sum_{k=1}^3 \widehat{\mu}(P_{cr}, \Xi_{n^*,k}) - \gamma_{\tau(k)} \widehat{\mu}(\tilde{P}_{cr}, \Xi_{n^*,k}) + \gamma_{\tau(k)} \widehat{\mu}(\tilde{P}_{cr}, \Xi_{m^*,k}), \quad (20)$$

$$\begin{aligned} \widehat{\sigma}^2 [\widehat{\mu}^{(CV+S)}] &= \frac{1}{3^2} \sum_{k=1}^3 \widehat{\sigma}^2 [\widehat{\mu}(P_{cr}, \Xi_{n^*,k})] - 2\gamma_{\tau(k)}^* \widehat{\delta} [\widehat{\mu}(P_{cr}, \Xi_{n^*,k}), \widehat{\mu}(\tilde{P}_{cr}, \Xi_{n^*,k})] + \\ &\quad (\gamma_{\tau(k)}^*)^2 \widehat{\sigma}^2 [\widehat{\mu}(\tilde{P}_{cr}, \Xi_{n^*,k})] + (\gamma_{\tau(k)}^*)^2 \widehat{\sigma}^2 [\widehat{\mu}(\tilde{P}_{cr}, \Xi_{m^*,k})], \end{aligned} \quad (21)$$

$$\gamma_{\tau(k)}^* = \frac{\widehat{\delta} [\widehat{\mu}(P_{cr}, \Xi_{n^*,\tau(k)}), \widehat{\mu}(\tilde{P}_{cr}, \Xi_{n^*,\tau(k)})]}{\widehat{\sigma}^2 [\widehat{\mu}(\tilde{P}_{cr}, \Xi_{n^*,\tau(k)})] + \widehat{\sigma}^2 [\widehat{\mu}(\tilde{P}_{cr}, \Xi_{m^*,\tau(k)})]}, \quad k = 1, 2, 3. \quad (22)$$

186 The structure of Eqs. (20), (21) and (22) reveals the essence of the Splitting approach. That is, when applying Control Variates to
 187 calculate either the estimator of the sought statistic (see Eq. (20)) or its variance (see Eq. (21)), one considers the k -th subset of
 188 samples for calculating $\widehat{\mu}$ while the optimal control parameter is calculated using the $\tau(k)$ -th subset of samples. As $k \neq \tau(k)$,
 189 such strategy effectively ensures that bias is avoided, as demonstrated in detail in [3]. Furthermore, it is noted from Eqs. (20),
 190 (21) and (22) that the implementation of the Splitting strategy does not demand any additional buckling analyses. Instead, it
 191 demands performing calculations of the different estimators and optimal control parameters over different subsets, which is quite
 192 cheap from a numerical viewpoint. In summary, the Splitting technique offers a convenient way to avoid the undesirable effects
 193 of bias while not increasing numerical costs.

194 4.3 | Estimation of the mean

195 The preceding section illustrates how a statistic of interest μ is estimated by means of Control Variates with Splitting, where μ
 196 can represent either mean or variance. When the focus is on estimating the mean of the non-linear buckling load (denoted as
 197 $\widehat{\mu}_1^{(CV+S)}$) by means of Control Variates with Splitting as well as the variance of that estimator (that is, $\widehat{\sigma}^2 [\widehat{\mu}_1^{(CV+S)}]$), Eqs. (20),
 198 (21) and (22) adopt the following specific form, respectively

$$\widehat{\mu}_1^{(CV+S)} = \frac{1}{3} \sum_{k=1}^3 \widehat{\mu}_1'(P_{cr}, \Xi_{n^*,k}) - \alpha_{\tau(k)} \widehat{\mu}_1'(\tilde{P}_{cr}, \Xi_{n^*,k}) + \alpha_{\tau(k)} \widehat{\mu}_1'(\tilde{P}_{cr}, \Xi_{m^*,k}), \quad (23)$$

$$\begin{aligned} \widehat{\sigma}^2 [\widehat{\mu}_1^{(CV+S)}] &= \frac{1}{3^2} \sum_{k=1}^3 \frac{\widehat{\mu}_{2,0}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*,k})}{n^*} - 2\alpha_{\tau(k)}^* \frac{\widehat{\mu}_{1,1}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*,k})}{n^*} \\ &\quad + (\alpha_{\tau(k)}^*)^2 \frac{\widehat{\mu}_{0,2}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*,k})}{n^*} + (\alpha_{\tau(k)}^*)^2 \frac{\widehat{\mu}_{0,2}(P_{cr}, \tilde{P}_{cr}, \Xi_{m^*,k})}{m^*}, \end{aligned} \quad (24)$$

$$\alpha_{\tau(k)}^* = \frac{\widehat{\mu}_{1,1}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, \tau(k)})}{\frac{\widehat{\mu}_{0,2}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, \tau(k)})}{n^*} + \frac{\widehat{\mu}_{0,2}(P_{cr}, \tilde{P}_{cr}, \Xi_{m^*, \tau(k)})}{m^*}}, \quad k = 1, 2, 3. \quad (25)$$

199 In the above equations, $\widehat{\mu}_1^*$ refers to the estimator of the mean (see Eq. (13)); $\alpha_{\tau(k)}^*$ represents the optimal control parameter
 200 associated the $\tau(k)$ -th subset which is related with the estimation of the mean; and the term $\widehat{\mu}_{p,q}$ represents the estimator of the
 201 bivariate central co-moment of order (p, q) between the non-linear and linear buckling loads. Detailed expressions for evaluating
 202 these co-moments are listed in A.

203 4.4 | Estimation of the variance

204 When the objective is to estimate the variance of the buckling load (denoted as $\widehat{\mu}_2^{(CV+S)}$) as well as the variance of that estimator
 205 (represented as $\widehat{\sigma}^2[\widehat{\mu}_2^{(CV+S)}]$) by means of Control Variates with Splitting, Eqs. (20), (21) and (22) must be formulated as

$$\widehat{\mu}_2^{(CV+S)} = \frac{1}{3} \sum_{k=1}^3 \widehat{\mu}_2(P_{cr}, \Xi_{n^*, k}) - \beta_{\tau(k)} \widehat{\mu}_2(\tilde{P}_{cr}, \Xi_{n^*, k}) + \beta_{\tau(k)} \widehat{\mu}_2(\tilde{P}_{cr}, \Xi_{m^*, k}), \quad (26)$$

$$\begin{aligned} \widehat{\sigma}^2[\widehat{\mu}_2^{(CV+S)}] &= \frac{1}{32} \sum_{k=1}^3 B_1(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}) - 2\beta_{\tau(k)} B_2(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}) \\ &+ (\beta_{\tau(k)})^2 B_3(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}) + (\beta_{\tau(k)})^2 B_4(P_{cr}, \tilde{P}_{cr}, \Xi_{m^*, k}), \end{aligned} \quad (27)$$

$$\beta_{\tau(k)}^* = \frac{B_2(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, \tau(k)})}{B_3(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, \tau(k)}) + B_4(P_{cr}, \tilde{P}_{cr}, \Xi_{m^*, \tau(k)})}, \quad k = 1, 2, 3. \quad (28)$$

206 In the last three equations, $\widehat{\mu}_2$ represents the estimator of the variance (see Eq. (14)); $\beta_{\tau(k)}^*$ is the optimal control parameter
 207 associated with the $\tau(k)$ -th subset which is related with the estimation of the variance; and B_1, B_2, B_3 and B_4 are real constants
 208 that are calculated by:

$$B_1(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}) = \frac{\widehat{\mu}_{4,0}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, j})}{n^*} - \frac{(n^* - 3)}{(n^* - 1)n^*} \widehat{\mu}_{2,0}^2(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}), \quad (29)$$

$$B_2(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}) = \frac{2\widehat{\mu}_{1,1}^2(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k})}{(n^* - 1)n^*} + \frac{\widehat{\mu}_{2,2}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k})}{n^*} - \frac{\widehat{\mu}_{2,0}\widehat{\mu}_{0,2}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k})}{n^*}, \quad (30)$$

$$B_3(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}) = \frac{\widehat{\mu}_{0,4}(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, j})}{n^*} - \frac{(n^* - 3)}{(n^* - 1)n^*} \widehat{\mu}_{0,2}^2(P_{cr}, \tilde{P}_{cr}, \Xi_{n^*, k}), \quad (31)$$

$$B_4(P_{cr}, \tilde{P}_{cr}, \Xi_{m^*, k}) = \frac{\widehat{\mu}_{0,4}(P_{cr}, \tilde{P}_{cr}, \Xi_{m^*, k})}{m^*} - \frac{(m^* - 3)}{(m^* - 1)m^*} \widehat{\mu}_{0,2}^2(P_{cr}, \tilde{P}_{cr}, \Xi_{m^*, k}). \quad (32)$$

209 Expressions for evaluating the co-moments associated with the calculation of constants B_1, B_2, B_3 and B_4 can be found in A.

210 5 | NUMERICAL EXAMPLES

211 Three numerical examples are presented using geometric non-linear quadrilateral shell elements with moderate rotations, as
 212 described in [42]. This four-node element is based on the isoparametric concept with linear shape functions. To prevent shear
 213 locking, the assumed natural strain (ANS) method is implemented. The element is incorporated into an extended version of
 214 the general finite element analysis program (FEAP) [36]. An interface has been developed to enable FEAP to be called from
 215 MATLAB, allowing the calculation of buckling load solutions within the Monte Carlo loop.

216 The linear buckling analysis is performed using the subspace iteration method as described in [5, 25]. In [5], the recommended
 217 number of eigenvalues to be determined by iteration is given by $N = \min\{2n, n+8\}$, where n is the number of desired eigenvalues.
 218 The first eigenvalue indicates the stability point. Therefore, n is set to one to maximize performance of the Control Variates
 219 approach. To compare the speed-up factors between the Control Variates approach and the Monte Carlo simulation for the
 220 presented examples, all calculations were performed on the same workstation with the following setup: $2 \times$ CPUs Intel Xeon
 221 E5-2667 v4 8 cores @ 3.20 GHz, 128 GB RAM, Win 10 x64

222 5.1 | Folded Plate

223 The first example of a folded plate, as depicted in Fig. 3, is to show that the Control Variates method works successfully in the
 presence of two types of stability points: bifurcation points and a limit point (snap-through). A distributed load p is applied to

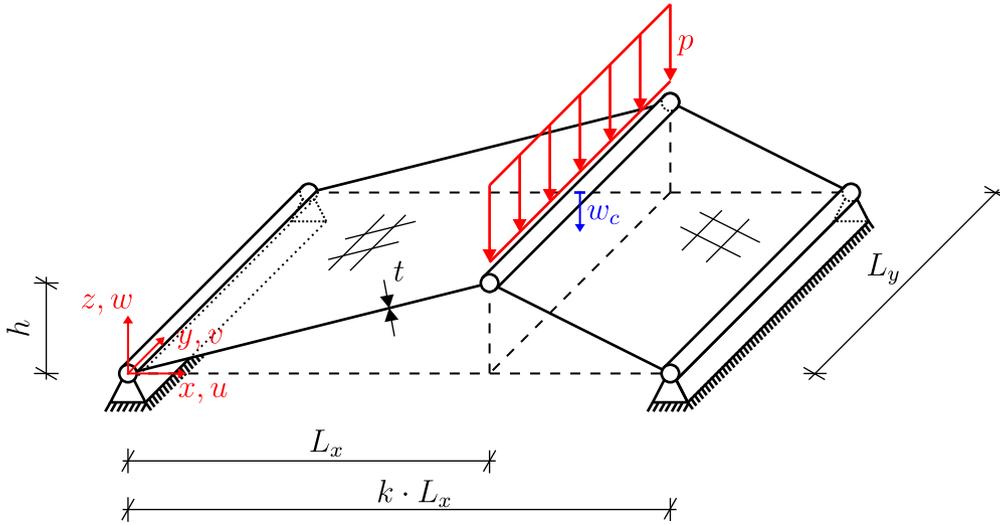


FIGURE 3 Folded Plate subjected to a distributed load

224 the plate at the hinged sleeve along the line $x = L_x$. The plate is simply supported at the edges $x = 0$ and $x = k \cdot L_x$, with boundary
 225 conditions $u = v = w = 0$. The lengths L_x and L_y are set to 100 cm and the plate thickness is defined as $t = 5$ cm. A length factor
 226 k allows to control the symmetry of the system and thus also the buckling behavior. For $k = 1.75$ and $k = 2.00$, the system
 227 is non-symmetric, whereas a symmetric system occurs when $k = 2.00$. For an initial analysis, the Young's modulus is set to
 228 $E = 1000$ kN/cm² and the height to $h = 10$ cm. Each subarea to the left and right of the hinged sleeve is discretized with 10×10
 229 elements. The vertical displacements of all nodes of the hinged sleeve are linked. The resulting load $P = p L_y$ is computed for an
 230 increasing vertical displacement w_c at $x = L_x, y = L_y/2$ using the arc-length method with a displacement control of $\Delta w = 0.02$ cm.
 231 The resulting load–displacement curve for a non-symmetric system ($k = 1.75$) is depicted in Fig. 4. Two bifurcation points
 232 (A) and (B) occur before the snap-through point (C). In addition, associated eigenvectors φ_{cr} are depicted in Fig. 4. The first
 233 eigenvectors in point (A) and (B) represent a local buckling of the left or right part of the folded plate, whereas the eigenvectors
 234 at point (C) indicate the snap-through as a global buckling failure mode. By applying the corresponding eigenvectors as small
 235 imperfections at points (A) and (B), the resulting secondary equilibrium paths, illustrated by the red dashed lines, can be traced.
 236 A non-linear buckling analysis yields a buckling load of $P_{cr} = 9.92$ kN at bifurcation point (A), $P_{cr} = 14.07$ kN at point (B) and
 237 $P_{cr} = 14.67$ kN at the snap-through point (C). The critical load from the linear buckling analysis is $\tilde{P}_{cr} = 10.83$ kN.

In the following, the length factor k , height h , and Young's modulus E are modeled as Gaussian random variables. Their mean
 values μ and standard deviations σ are listed in Table 2. The stochastic non-linear buckling response is analyzed via a Monte
 Carlo simulation with 5000 samples, resulting in a mean value of the critical buckling load $\mu_{P_{cr}} = 9.94$ kN and a coefficient
 of variation of $\delta_{P_{cr}} = 18.82\%$. All buckling loads correspond to the load, where the first diagonal element becomes negative.
 Despite the significant influence of the three random input parameters defined in Table 2 on the buckling behavior, a strong
 correlation of $\rho = 0.99$ between linear and non-linear buckling analyses is evident, see Fig. 5. These conditions are well-suited

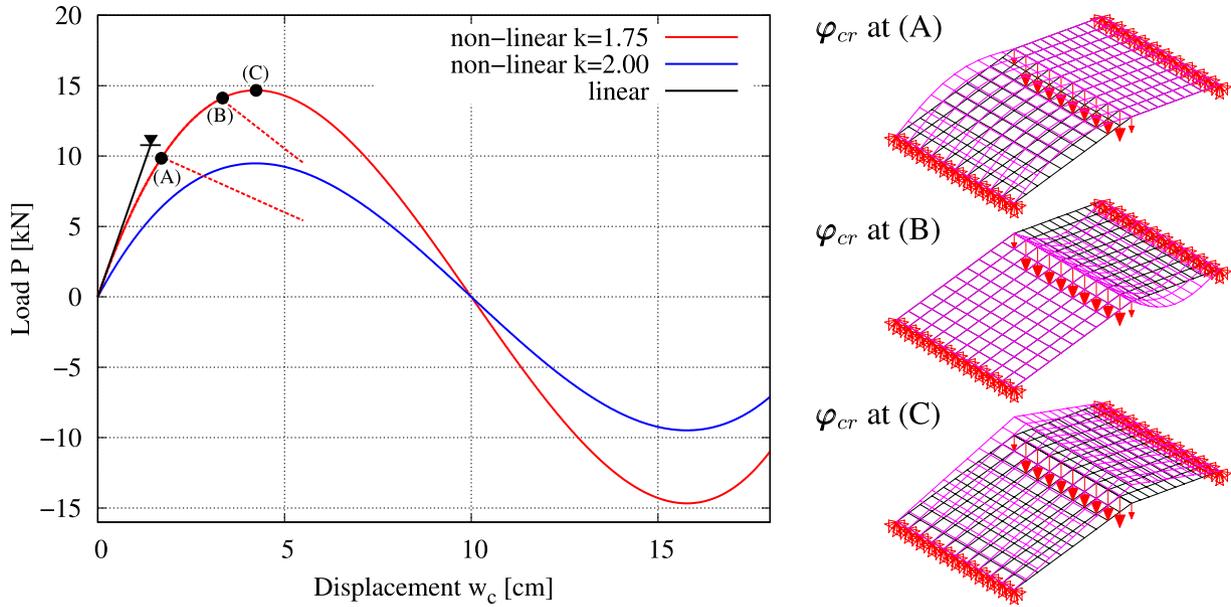


FIGURE 4 Load-displacement curves of the folded plate

TABLE 2 Folded plate: Quantification of material and geometric parameters as Gaussian random variables

parameter	mean value μ	standard deviation σ
length factor k [-]	1.75	0.05
height h [cm]	10	1
Young's modulus E [kN/cm ²]	1000	100

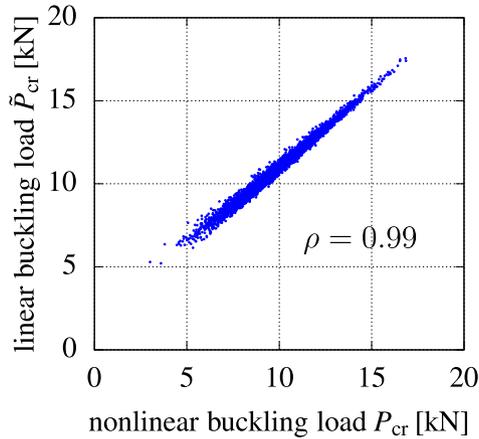


FIGURE 5 Correlation between buckling loads calculated using non-linear P_{cr} and linear analysis \tilde{P}_{cr}

for applying the Control Variates method. The results are given in Table 3. The estimates obtained using Monte Carlo and Control Variates with Splitting are nearly identical, as indicated by the coefficient of variation δ in Table 3. To compare the numerical efforts between the two approaches for estimating second-order statistics, the total number of analyses performed for Control Variates is expressed as an equivalent number of analyses n_e (relative to plain Monte Carlo simulation). This equivalent number of analyses is calculated by

$$n_e = n + \frac{n + m}{f_s}, \quad (33)$$

TABLE 3 Estimates of second-order statistics for the buckling load of the folded plate.

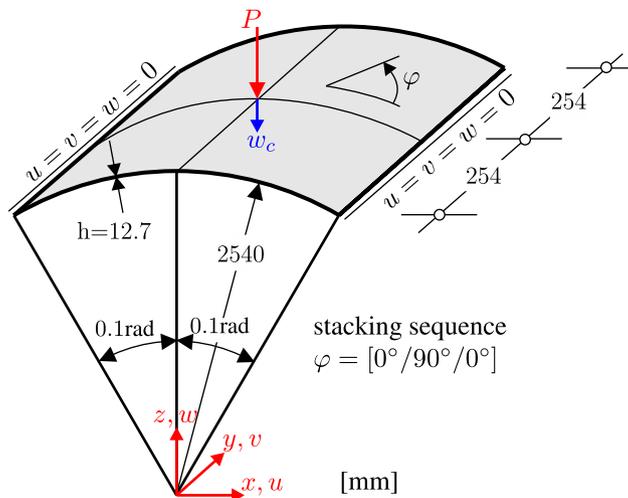
Approach	Monte Carlo	Control Variates with Splitting
n	1750	150
m	–	3000
n_e	1750	736
$\widehat{\mu}_1 N$	9.99	9.97
$V[\widehat{\mu}_1] N^2$	0.002	0.002
$\delta_{\widehat{\mu}_1}$	0.45%	0.39%
$\widehat{\mu}_2 N^2$	3.582	3.614
$V[\widehat{\mu}_2] N^4$	0.015	0.018
$\delta_{\widehat{\mu}_2}$	3.45%	3.65%
$\delta_{P_{cr}}$	19%	19%

where n is the number of samples for the non-linear buckling analysis and m is the number of samples for the linear buckling analysis. The speedup factor f_s is defined as the ratio between the execution time of one non-linear buckling analysis and one linear buckling analysis. The number of $n_e = 736$ taken from Table 3 means that the Control Variates requires only a number of 736 equivalent simulations instead of 1750 Monte Carlo simulations. In other words, for a sample size of $n = 150$ and $m = 3000$, a speedup factor of $f_s = 5.38$ is obtained, indicating that the linear buckling analysis is 5.38 times more computationally efficient than the non-linear buckling analysis.

Finally, it should be noted that even if a bifurcation path exists in the pre-buckling stage, the method still performs successfully. The only essential requirement is a strong correlation between the linear and non-linear buckling analyses, while the type of stability point plays a subordinate role. If a sufficiently strong correlation exists, the Control Variates method yields reliable results.

5.2 | Composite shell panel

This example demonstrates the use of the Control Variates approach to predict buckling loads in a composite shell panel subjected to a single load and random geometric imperfections. For details on the shell panel model, see [41]. Fig. 6 illustrates the system modeled with a 30×30 FE mesh. The panel is simply supported along the two lateral edges. System symmetry cannot be exploited due to applied random geometrical imperfections.

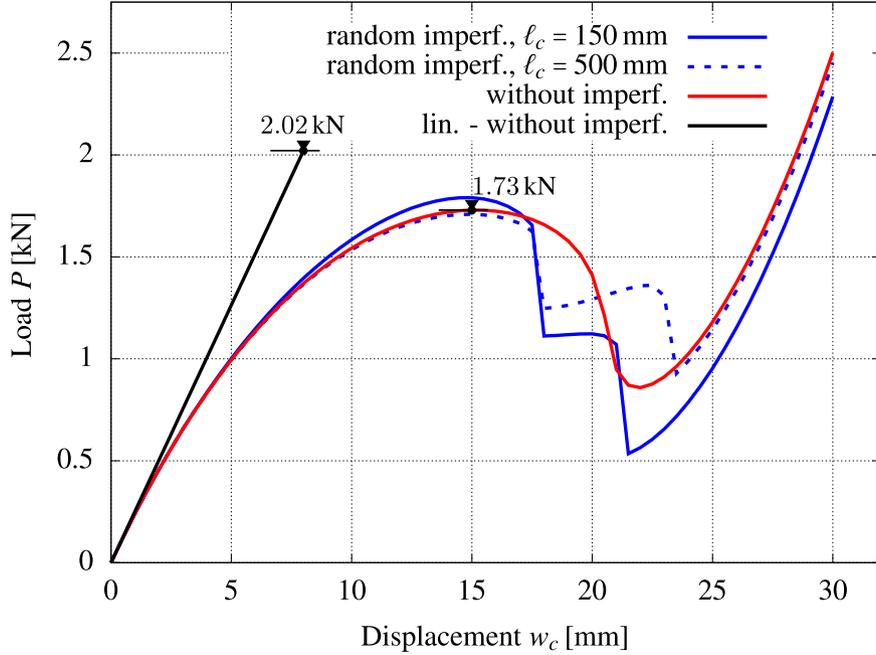
**FIGURE 6** Cylindrical composite shell panel subjected to a single load

The transversely isotropic material properties of the laminate are provided in Table 4. The panel consists of a three-layer laminate with a total thickness of $h = 12.7$ mm. The corresponding stacking sequence is given in Fig. 6.

TABLE 4 Transversal isotropic material properties of the laminate

E_{11} [N/mm ²]	E_{22} [N/mm ²]	G_{12} [N/mm ²]	G_{13} [N/mm ²]	G_{23} [N/mm ²]	ν_{12} [-]
3300	1100	660	660	450	0.3

256 First of all, the buckling behavior is analyzed without any imperfections. The load-displacement curves, evaluating the load
 257 P versus the vertical displacement w_c at the center of the panel, are depicted in Fig. 7. These curves are obtained using the
 258 arc-length method with a displacement control $\Delta w = 0.5$ mm.

**FIGURE 7** Load-displacement curves of the composite shell panel

259 A pronounced non-linear pre-buckling behavior is observed. Consequently, the non-linear buckling analysis using the criterion
 260 in Eq. (12) results in a critical load of $P_{cr} = 1.73$ kN, while the linear buckling analysis according to Eq. (9) yields a higher
 261 critical load of $\tilde{P}_{cr} = 2.02$ kN.

262 The next step is to generate random geometric imperfections of the cylindrical shell shape (reference surface) modeled as
 263 Gaussian random fields using the Karhunen-Loève Expansion (KLE). For the fundamentals to random field modeling, see, e.g.,
 264 [35, 37]. Accordingly, the random geometric deviations in radial direction of the shell panel can be expressed as

$$\hat{w}_{\text{rad}}(\mathbf{x}, \theta) = \mu + \sum_{i=1}^M \sqrt{\lambda_i} \xi_i(\theta) \varphi_i(\mathbf{x}). \quad (34)$$

265 Here, the mean value μ is set to zero to model geometric imperfections that vary around the reference surface. The parameter
 266 $\xi_i(\theta)$ is a standard normal distributed random variable, $\varphi_i(\mathbf{x})$ are the eigenfunctions and λ_i the eigenvalues of the covariance
 267 matrix formulated for the FE mesh with M nodes. This covariance matrix is assumed to be homogeneous

$$C(\tau) = \sigma^2 \rho(\tau), \quad (35)$$

where the variance σ^2 is set to one in this example. The autocorrelation function (acf) $\rho(\tau)$ is defined as a function of the relative distance τ of two FE nodes $\mathbf{x}_i, \mathbf{x}_j$

$$\rho(\mathbf{x}_i, \mathbf{x}_j) = \rho(\tau) \quad \text{with} \quad \tau = \mathbf{x}_j - \mathbf{x}_i. \quad (36)$$

In practical applications, analytical models for autocorrelation are frequently used [9, 13]. For the presented example, the Whittle–Matérn acf is chosen to generate the shell imperfections

$$\rho(\tau) = \frac{2^{1-\nu}}{\Gamma(\nu)} \left(\sqrt{2\nu} \frac{\tau}{\ell_c} \right)^\nu K_\nu \left(\sqrt{2\nu} \frac{\tau}{\ell_c} \right), \quad (37)$$

where K_ν is the modified Bessel function of the second kind and $\Gamma(\nu)$ denotes the Gamma function. The parameter ℓ_c is the correlation length and ν is the so-called “smoothness” parameter, which is set to $\nu = 1.5$.

The correlation length influences the imperfection shape, thereby it also influences the buckling behavior and the second-order statistics of the buckling load. Fig. 8 illustrates the mean $\mu_{P_{cr}}$ and the coefficient of variation $\delta_{P_{cr}}$ of the buckling load as functions of the correlation length.

For the small correlation length $\ell_c = 150$ mm, a minimal mean value $\mu_{P_{cr}} = 1.68$ kN and a maximum coefficient of variation $\delta_{P_{cr}} = \sigma_{P_{cr}}/\mu_{P_{cr}} = 10\%$ can be observed. From this point, the mean value converges to the imperfection-free buckling load $P_{cr} = 1.73$ kN and the coefficient of variation decreases. This is because, for large correlation lengths, the effect of geometric imperfections vanish. Examining the vertical axis of Fig. 8 (left), the mean value shows small variation. However, the correlation length has a significant influence on the coefficient of variations, see Fig. 8 (right). Therefore, two different correlation lengths are selected to investigate the effectiveness of Control Variates: the smaller length of $\ell_c = 150$ mm, which results in wavy realizations, and a larger length of $\ell_c = 500$ mm, which produces more uniform imperfection shapes as depicted in Fig. 9.

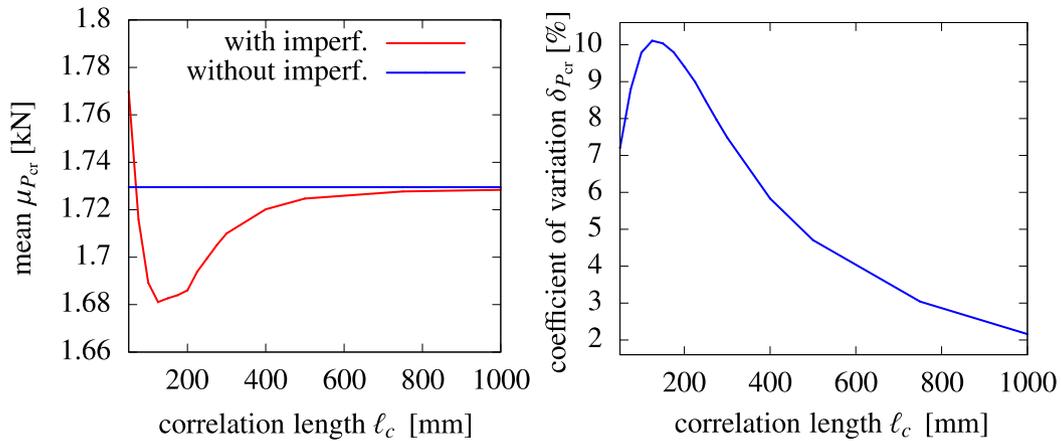


FIGURE 8 Mean $\mu_{P_{cr}}$ (left) and coefficient of variation $\delta_{P_{cr}}$ (right) of the buckling load as a function of the correlation length

In addition, the load-displacement curves for two realizations of the selected correlation lengths are depicted Fig. 7. For these realizations, a small difference in the buckling load can be observed, consistent with the slight deviations in the mean value shown in Fig. 8 (left). In Fig. 10, the eigenvectors φ_{cr} at the stability point are depicted for the panel without and with the imperfection from Fig. 9 (left).

The eigenvectors are quite similar. However, the eigenvector with imperfections is non-symmetric, leading to different post-buckling behavior. Table 5 summarizes the second-order statistics for both correlation lengths, as evaluated using a Monte Carlo simulation with 5000 realizations.

For instance, Fig. 11 shows the Monte Carlo convergence test for generated random imperfections with the correlation length $\ell_c = 150$ mm. For 1000 samples, the relative error of $\delta_{P_{cr}}$ compared to the reference solution provided in Table 5 is smaller than 5%.

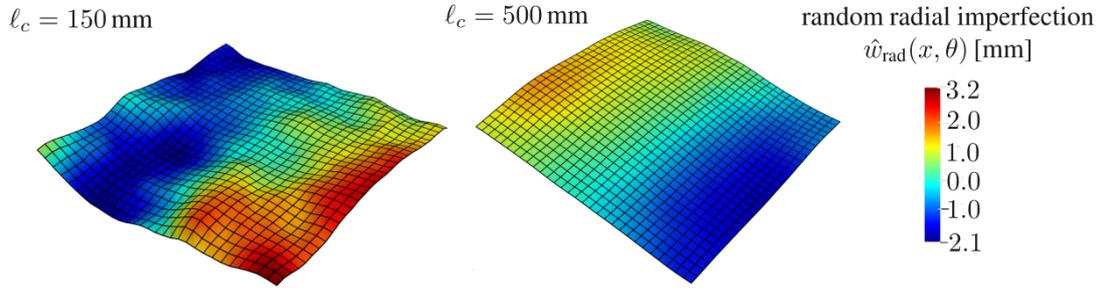


FIGURE 9 Random geometrical imperfection with respect to the shell surface for $\ell_c = 150$ mm (left) and $\ell_c = 500$ mm (right)

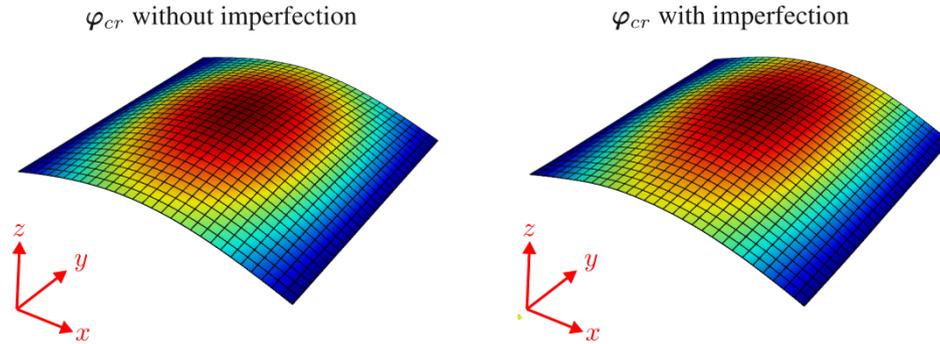


FIGURE 10 Eigenvectors φ_{cr} at the stability point of the shell panel without (left) and realization with random imperfection for the correlation length $\ell_c = 150$ mm (right)

TABLE 5 Second-order statistics of the buckling load for different correlation lengths

correlation length ℓ_c [mm]	$\mu_{P_{cr}}$ [kN]	$\delta_{P_{cr}}$ [%]
150	1.68	10.0
500	1.72	4.7

290 Fig. 12 shows the correlations between buckling loads calculated using non-linear (P_{cr}) and linear analysis (\tilde{P}_{cr}) for the small
 291 and large correlation length. For each correlation length, 5000 random fields are generated and the buckling loads with the linear
 292 and non-linear buckling analysis are calculated.

293 It is noteworthy that a strong correlation of $\rho = 0.97$ is observed even for the small correlation length of $\ell_c = 150$ mm.
 294 Compared to the larger length $\ell_c = 500$ mm, the correlation between linear and non-linear buckling analyses decreases only
 295 slightly. This is a good condition for using Control Variates. To calculate the linear solutions is 4.8 faster than performing the
 296 non-linear buckling analyses. However, for the small correlation length, the subspace eigenvalue solver requires sometimes more
 297 iteration steps to compute the correct eigenvalue. This means that the chosen eigenvalue solver and its properties significantly
 298 influence the effectiveness of the Control Variates approach.

299 The results of the estimates of second-order statistics using the Control Variates approach are provided for both correlation
 300 lengths in Table 6. The estimates produced using Monte Carlo and Control Variates with Splitting are practically identical in
 301 terms of both accuracy and precision, measured in terms the coefficient of variation δ in Table 6. For a smaller correlation length
 302 $\ell_c = 150$ mm, the number of $n_e = 285$ according to Eq. (33) means that the Control Variates requires only a number of 285
 303 equivalent simulations instead of 600 Monte Carlo simulations. In other words, for the number of samples $n = 90$ and $m = 999$, a
 304 speedup factor of $f_s = 5.58$ is achieved, which means that the linear buckling analysis is 5.58 times faster than the non-linear
 305 buckling analysis.

306 For the larger correlation length, Control Variates performs more efficiently with $f_s = 5.75$ and an equivalent number of 227
 307 analyses instead of 600. This can be explained by analyzing Fig. 12, as the correlation coefficient between linear and non-linear
 308 buckling loads is $\rho = 0.99$, which is larger than the correlation coefficient associated with a shorter correlation length. In other

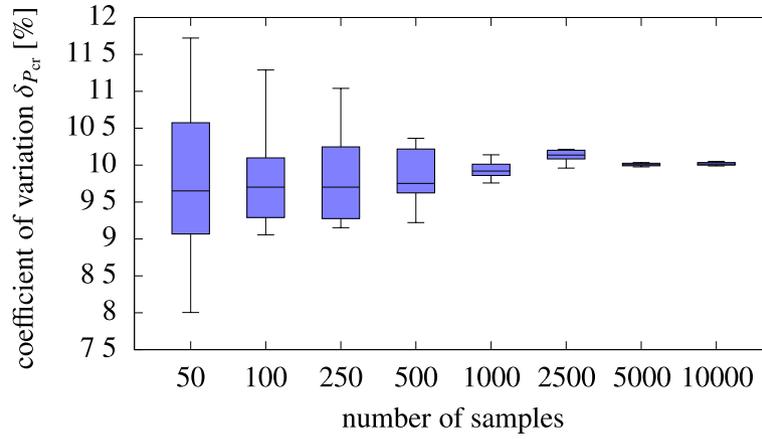


FIGURE 11 Monte Carlo convergence test for correlation length $\ell_c = 150$ mm

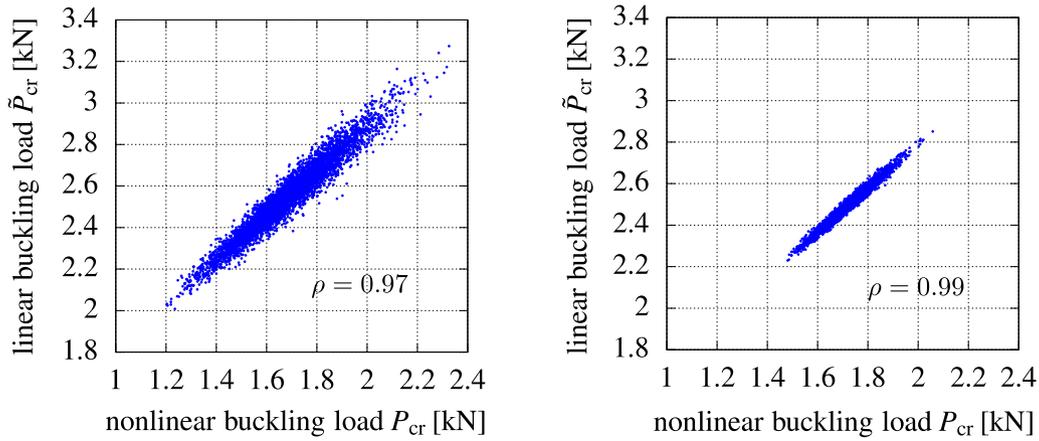


FIGURE 12 Comparison between buckling loads calculated using non-linear (P_{cr}) and linear analysis (\tilde{P}_{cr}) for the correlation lengths $\ell_c = 150$ mm (left) and $\ell_c = 500$ mm (right)

309 words, for the case of a longer correlation length, linear buckling analysis provides a better approximation for calculating the
 310 exact buckling load. In addition, the subspace eigenvalue solver converges faster for smooth imperfection shapes.

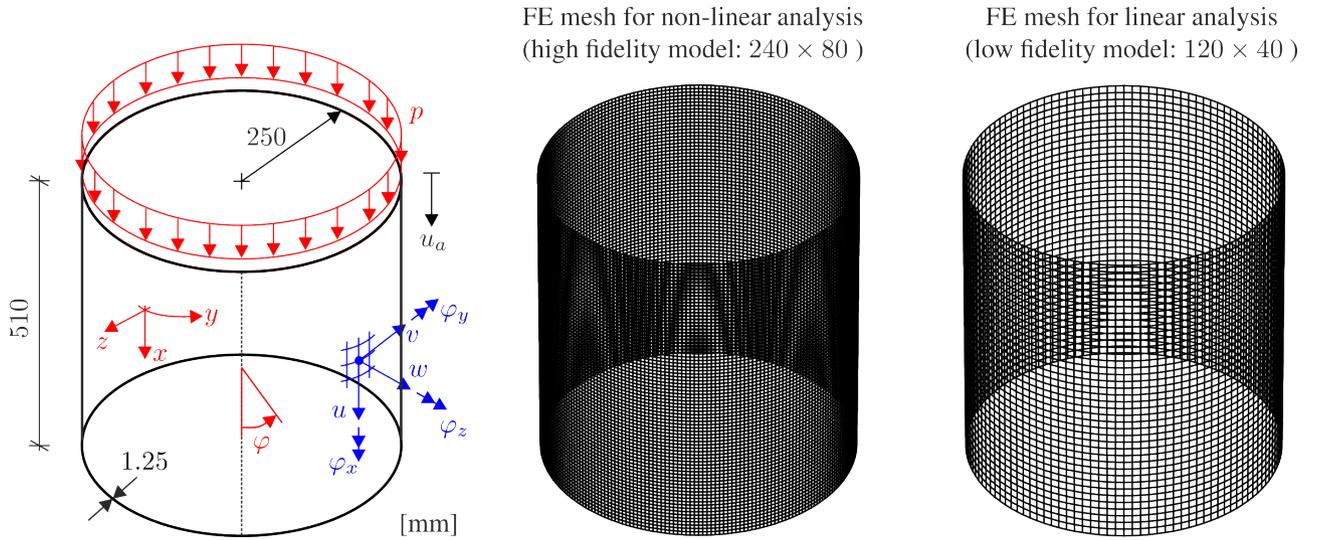
311 When examining the results of Table 6, it is observed that while the mean buckling load increases slightly with the correlation
 312 length, the variance decreases considerably. This is a very interesting behavior, as in problems of linear static stochastic FE
 313 analysis, usually the opposite behavior is observed [45, 10]. This highlights the non-linear nature of the problem at hand.

314 5.3 | Composite cylinder

315 Buckling analysis of cylindrical shells is particularly challenging from both theoretical and numerical perspectives, and it is
 316 associated with high computational costs. Therefore, the effectiveness of the Control Variates method is investigated for an
 317 imperfection-sensitive composite cylinder subjected to random geometric imperfections. For this purpose, cylinder Z23, derived
 318 from [28, 39], is analyzed. The corresponding FE model is depicted in Fig. 15. The cylinder has a length of $L = 510$ mm and a
 319 radius of $R = 250$ mm, resulting in a circumference approximately three times larger than its length. Based on a convergence
 320 study, an FE mesh with 240 shell elements in the circumferential direction and 80 elements in the axial direction is chosen.
 321 This approximately regular FE mesh effectively captures the critical buckling modes and is used as the high fidelity model for
 322 the non-linear buckling analysis. In contrast, a coarser mesh of 120×40 is used as a low-fidelity model to perform the linear
 323 buckling analysis within the Control Variates approach. The cylinder's laminate consists of 10 layers with a stacking sequence
 324 of $[\pm 60^\circ, 0_2^\circ, \pm 68^\circ, \pm 52^\circ, \pm 37^\circ]$, corresponding to the fiber orientation φ depicted in Fig. 15. A single layer has a thickness

TABLE 6 Estimates of second-order statistics for the buckling load of the composite shell panel.

Approach	Monte Carlo	Control Variates with Splitting	Monte Carlo	Control Variates with Splitting
ℓ_c [mm]		150		500
n	600	90	600	60
m	–	999	–	900
n_e	600	285	600	227
$\widehat{\mu}_1$ [kN]	1.7	1.69	1.72	1.73
$V[\widehat{\mu}_1]$ [kN ²]	5×10^{-5}	4.1×10^{-5}	1.0×10^{-5}	8.1×10^{-6}
$\delta_{\widehat{\mu}_1}$	0.4%	0.4%	0.2%	0.2%
$\widehat{\mu}_2$ [kN ²]	0.03	0.03	0.006	0.006
$V[\widehat{\mu}_2]$ [kN ⁴]	2.8×10^{-6}	2.8×10^{-6}	1.2×10^{-7}	1.0×10^{-7}
$\delta_{\widehat{\mu}_2}$	5.6%	5.8%	5.6%	5.6%
$\delta_{P_{cr}}$	10.2%	10%	4.6%	4.4%

**FIGURE 13** FE model of the composite cylinder

of 0.125 mm, leading to a total shell thickness of $t = 1.25$ mm. The material parameters are given in Table 7. The cylinder is

E_{11} [N/mm ²]	E_{22} [N/mm ²]	G_{12} [N/mm ²]	G_{23} [N/mm ²]	ν_{12} [–]
123 550	8 708	5 695	3 400	0.319

TABLE 7 Material parameters for the composite cylinder Z23

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clamped at both ends. At the lower edge, the boundary conditions are $u = v = w = 0, \varphi_x = \varphi_y = 0$. Only vertical displacements of the nodes at the upper edge are allowed $u = u_a$, while $v = w = 0, \varphi_x = \varphi_y = 0$ remain fixed. The resulting axial load is computed as $P = 2\pi R p$ for an increasing vertical displacement u_a using the arc-length method. The displacement step size significantly influences the computational time. Therefore, five initial coarse displacement steps of $\Delta u_a = 0.1$ mm are applied. This coarse steps are estimated based on the lowest expected buckling load of the cylinder under random imperfections. Subsequently, the cylinder is further loaded with a smaller displacement step size of $\Delta u_a = 0.01$ mm until the first zero diagonal element D_{ii} occurs in the tangent stiffness matrix.

In this example, the random geometric imperfections are generated using the EOLE (Expansion Optimal Linear Estimation) method from [27]. The method allows to represent the random field with only a few random variables by minimizing the variance error. The main advantage is that the covariance matrix is only required on a sub-set of field nodes, the so-called "random field mesh". Thus, a coarser mesh can be defined for the random field compared to the finer mesh required for the FE analysis. The

expression for the series to compute random radial imperfections of the cylinder is

$$\hat{w}_{\text{rad}}(\mathbf{x}, \theta) = \mu + \left(\sum_{i=1}^M \frac{\xi_i(\theta)}{\sqrt{\lambda_i}} \varphi_i(\mathbf{x}^S) \right) C(\mathbf{x}^S, \mathbf{x}), \quad (38)$$

with the standard normal distributed random variable $\xi_i(\theta)$, the vector $\mathbf{x}^S = [\mathbf{x}_1^S \dots \mathbf{x}_i^S \dots \mathbf{x}_M^S]$ of M random field nodes and the vector $\mathbf{x} = [\mathbf{x}_1 \dots \mathbf{x}_j \dots \mathbf{x}_N]$ of N nodes in the full domain (e.g. FE nodes). Consequently, $C(\mathbf{x}^S, \mathbf{x})$ denotes the covariance matrix, which contains the covariances between random field nodes and FE nodes. The eigenfunctions $\varphi_i(\mathbf{x}^S)$ and eigenvalues λ_i are obtained from the covariance matrix $C(\mathbf{x}_i^S, \mathbf{x}_j^S)$ based on the random field mesh. Both covariance matrices $C(\mathbf{x}^S, \mathbf{x})$ and $C(\mathbf{x}^S, \mathbf{x}^S)$ are calculated using the homogeneous correlation structure in Eq. (36) and the Whittle–Matérn acf in Eq. (37). Depending on the smoothness of the random field, the number of random field nodes can be smaller than the number of FE nodes. This allows to reduce the size of the eigenvalue problem of the covariance matrix. In this example, a random field mesh with $M = 60 \times 20$ nodes is defined based on a convergence study, where the second-order statistics of the buckling load are evaluated. In Eq. (38), the constant mean value μ is set to zero and the standard deviation of the field is chosen as $\sigma = 1$. Considering the presented measured geometric imperfections of the cylinder in [39], the random imperfections are scaled to a peak to peak value of 2 mm.

In a first study, the stochastic buckling behavior is analyzed for different correlation length. The results of the mean $\mu_{P_{\text{cr}}}$ and the coefficient of variation $\delta_{P_{\text{cr}}}$ of the buckling load as functions of the correlation length are depicted in Fig. 14. For each

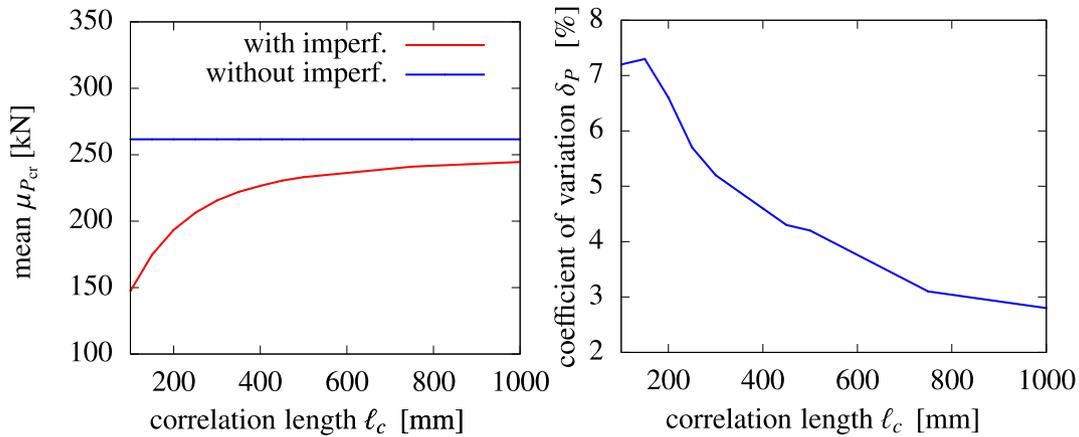


FIGURE 14 Mean $\mu_{P_{\text{cr}}}$ (left) and coefficient of variation $\delta_{P_{\text{cr}}}$ (right) of the buckling load as a function of the correlation length

correlation length, a Monte Carlo simulation with 500 samples is performed. In Fig. 14 (left), the resulting nonlinear buckling load of $P_{\text{cr}} = 261.64$ kN of the cylinder without imperfections is depicted as a horizontal line. Due to the applied scaling of the random geometric imperfections, the mean of the buckling load asymptotically approaches a value below the nonlinear buckling load as the correlation length increases. In contrast, the coefficient of variation decreases with increasing correlation length, see Fig. 14 (right). A maximum coefficient of variation of $\delta_{P_{\text{cr}}} = \sigma_{P_{\text{cr}}} / \mu_{P_{\text{cr}}} = 7.3\%$ results for the correlation length of $\ell_c = 150$ mm, which is selected for the following investigations. With respect to the mean of the buckling load $\mu_{P_{\text{cr}}} = 174.70$ for the chosen correlation length of $\ell_c = 150$ mm, the knockdown factor (KDF) is $\mu_{P_{\text{cr}}} / P_{\text{cr}} = 174.7 / 261.64 = 0.67$.

The load-displacement curves of the composite cylinder with and without random radial imperfections are depicted in Fig. 14. A typical linear pre-buckling behavior of a cylindrical shell can be observed. At the stability point (A) of the cylinder without imperfections, the initial post-buckling mode φ_{cr} is depicted. Due to the clamped edges, where the radial expansion of the cylinder is suppressed, the buckling mode is characterized by radial displacements at top and bottom of the cylinder. Furthermore, a sample of the scaled radial imperfection is depicted in Fig. 14, which is associated to the blue load-displacement curve and a buckling load of $P_{\text{cr}} = 165.86$ kN (KDF of 0.63).

The buckling load of the linear analysis using the high fidelity model (FE mesh with 240×80 elements) is $\tilde{P}_{\text{cr}} = 276.96$ kN. Whereas, the linear buckling load using the low fidelity model (FE mesh with 140×40 elements) is $\tilde{P}_{\text{cr}} = 301.02$ kN. The model behaves significantly stiffer. However, the correlation between the linear and nonlinear analysis decreases only slightly. This can be observed in the correlation plots in Fig. 16, where only the model of the linear analysis is changed. The nonlinear buckling

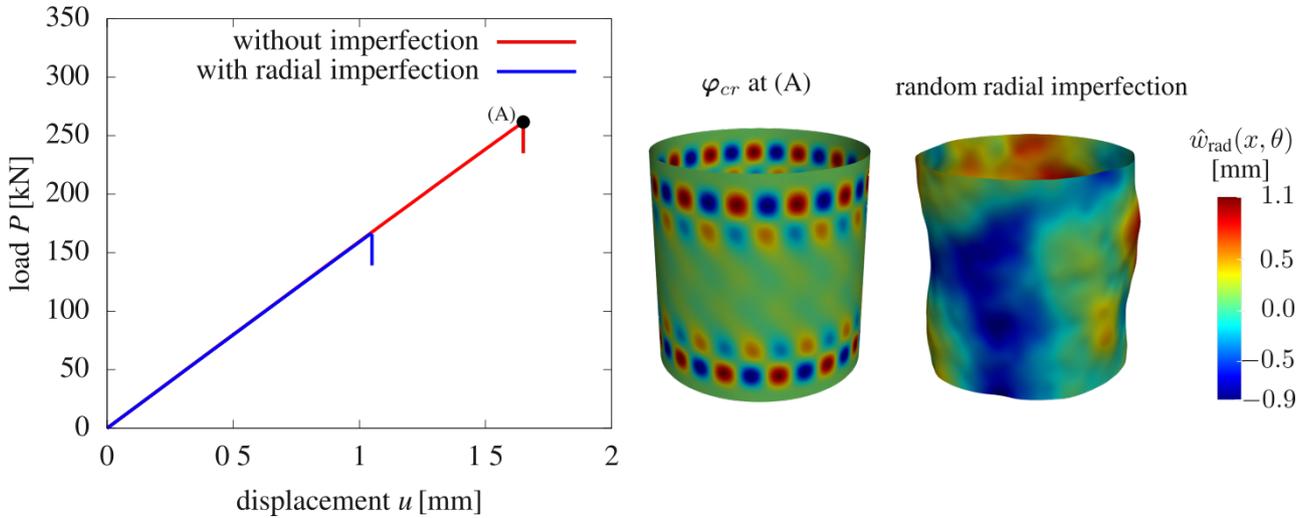


FIGURE 15 Load-displacement curve of the composite cylinder Z23 with the first eigenvector at the stability point of the cylinder without imperfection and a random radial imperfection, magnified (x20)

analysis is performed with the high fidelity model. On the given computational setup, the linear analysis using the low-fidelity

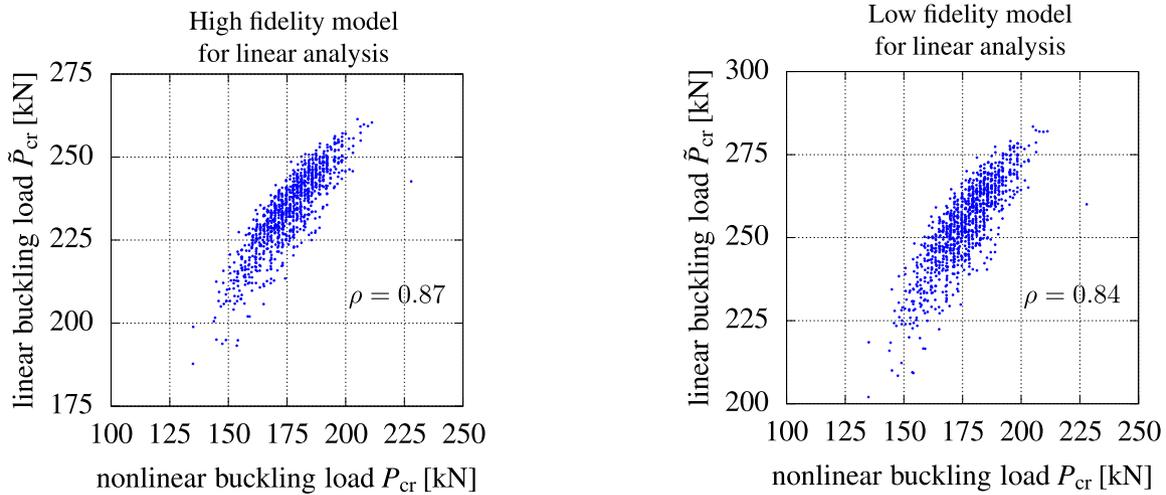


FIGURE 16 Comparison between buckling loads calculated using non-linear (P_{cr}) and linear analysis (\tilde{P}_{cr}) using the high fidelity model (left) and the low fidelity model (right) for the linear buckling analysis

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363 model is 13 times faster compared to the high-fidelity model. Thus, the effectiveness of the Control Variates method can be
 364 significantly enhanced by performing the linear buckling analysis using the low-fidelity model. The results of Control Variates
 365 approach compared to the Monte Carlo simulation are given in Table 8. According to Eq. (33), the Control Variates requires only
 366 a number of $n_e = 285$ equivalent simulations instead of 390 Monte Carlo simulations. This results in a speed-up factor of 20.69.
 367 Despite the correlation is approximately 10% lower than in the other examples, the computational time is significantly reduced.
 368 Thus, the CV method can serve as an effective tool for increasing the efficiency of stochastic analyses in shell buckling.

TABLE 8 Estimates of second-order statistics for the buckling load for the composite cylinder.

Approach	Monte Carlo	Control Variates with Splitting
n	390	210
m	–	990
n_e	390	268
$\widehat{\mu}_1 N$	176.5	176.1
$V[\widehat{\mu}_1] N^2$	0.41	0.33
$\delta_{\widehat{\mu}_1}$	0.4%	0.3%
$\widehat{\mu}_2 N^2$	161.4	154.2
$V[\widehat{\mu}_2] N^4$	137.8	129.8
$\delta_{\widehat{\mu}_2}$	7.3%	7.4%
$\delta_{P_{cr}}$	7.2%	7.1%

6 | CONCLUSIONS

The paper highlights a promising approach to estimate the second-order statistics of buckling loads. By integrating results from both linear and non-linear buckling analyses, this method achieves enhanced accuracy and significantly reduced computational costs compared to full-scale Monte Carlo simulations. However, the effectiveness of the Control Variates method for buckling problems depends on various factors and can be further optimized. One factor is the eigenvalue solver used for the linear buckling analysis. Therefore, other solvers have to be tested in this context.

Another important factor are the correlation properties of the random field, as random geometric imperfections. In this paper the influence of different correlation lengths are studied, but the correlation function and its differentiability may also influence the computational efficiency of calculating second-order statistics with Control Variates. Defining relevant correlation lengths and functions based on experimental data is essential for quantifying the benefits of Control Variates in realistic scenarios. To account for epistemic uncertainties, correlation parameters can be quantified using polymorphic uncertainty models. The application of the Control Variates method in the context of polymorphic uncertainties (imprecise probabilities) is also conceivable.

Finally, the composite shell panel discussed is primarily an academic example. The approach should be tested on a range of structures, including stiffened panels, fiber-steered composites, cylindrical shells, and large-scale structures, to assess its broader applicability. New ideas for further research can be summarized as follows:

- Efficient eigenvalue solvers for Control Variates
- Study to the influence of the correlation functions and their differentiability on estimating second-order statistics using Control Variates.
- Application of Control Variates for various structures, such as cylindrical shells, stiffened panels, fiber-steered composites and large-scale structures
- Application of Control Variates in the framework of polymorphic uncertainties (imprecise probabilities)
- Extension of the current framework from a single variable for performing Control Variates (in this case, buckling load from linear analysis) to several variables (for example, higher-order buckling loads from linear analysis). It is to be noted that extension towards several variables within the framework of Control Variates is possible, as discussed in e.g. [3].

ACKNOWLEDGMENTS

Financial support was provided by the Deutsche Forschungsgemeinschaft (DFG, German Research Foundation) in the framework of project 511267658. This support is gratefully acknowledged.

CONFLICT OF INTEREST

The authors declare no potential conflict of interests.

DATA AVAILABILITY STATEMENT

The data and materials that support the findings of this paper are available upon request to the corresponding author.

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478 APPENDIX

479 A BIVARIATE CENTRAL CO-MOMENTS

Bivariate central co-moments between the non-linear buckling load P_{cr} and linear buckling load \tilde{P}_{cr} are denoted as $\mu_{p,q}(P_{cr}, \tilde{P}_{cr})$, where the pair of integer numbers p and q represent the order associated with P_{cr} and \tilde{P}_{cr} , respectively. Monte Carlo simulation is employed to estimate these co-moments considering a sample set Ξ_l with l independent, identically distributed samples of Ξ . The list of co-moments required to implement the expressions of this paper are listed below. This list has been produced using the software package *mathStatica* [34]. Note that the expressions below consider the auxiliary variable $s_{p,q}$, which is defined as $s_{p,q} = \sum_{i=1}^l (P_{cr}(\xi^{(i)}))^p (\tilde{P}_{cr}(\xi^{(i)}))^q$, where $\xi^{(i)}$ is the i -th sample of the sample set Ξ_l .

$$\widehat{\mu_{1,1}} = \frac{ls_{1,1} - s_{0,1}s_{1,0}}{(l-1)l} \quad (A1)$$

$$\begin{aligned} \widehat{\mu_{2,2}} = & \frac{1}{(l-3)(l-2)(l-1)l} \left((-2l^2 + 4l - 6) s_{2,1}s_{0,1} + (-2l^2 + 4l - 6) s_{1,0}s_{1,2} + \right. \\ & (l^3 - 2l^2 + 3l) s_{2,2} + ls_{2,0}s_{0,1}^2 + 4ls_{1,0}s_{1,1}s_{0,1} + ls_{0,2}s_{1,0}^2 + \\ & \left. (6-4l)s_{1,1}^2 + (3-2l)s_{0,2}s_{2,0} - 3s_{1,0}^2s_{0,1}^2 \right) \end{aligned} \quad (A2)$$

$$\begin{aligned} \widehat{\mu_{4,0}} = & \frac{1}{(l-3)(l-2)(l-1)l} \left((-4l^2 + 8l - 12) s_{3,0}s_{1,0} + (l^3 - 2l^2 + 3l) s_{4,0} + \right. \\ & \left. 6ls_{2,0}s_{1,0}^2 + (9-6l)s_{2,0}^2 - 3s_{1,0}^4 \right) \end{aligned} \quad (A3)$$

$$\begin{aligned} \widehat{\mu_{0,4}} = & \frac{1}{(l-3)(l-2)(l-1)l} \left((-4l^2 + 8l - 12) s_{0,3}s_{0,1} + (l^3 - 2l^2 + 3l) s_{0,4} + \right. \\ & \left. 6ls_{0,2}s_{0,1}^2 + (9-6l)s_{0,2}^2 - 3s_{0,1}^4 \right) \end{aligned} \quad (A4)$$

Squared co-moments as well as co-moment products are estimated with the equations listed below [34].

$$\begin{aligned} \widehat{\mu_{1,1}^2} = & \frac{1}{(l-3)(l-2)(l-1)l} \left((l^2 - 3l + 2) s_{1,1}^2 + (l - l^2) s_{2,2} + \right. \\ & (2-2l)s_{1,0}s_{1,1}s_{0,1} + (2l-2)s_{2,1}s_{0,1} + (2l-2)s_{1,0}s_{1,2} + \\ & \left. s_{1,0}^2s_{0,1}^2 - s_{2,0}s_{0,1}^2 - s_{0,2}s_{1,0}^2 + s_{0,2}s_{2,0} \right) \end{aligned} \quad (A5)$$

$$\widehat{\mu_{2,0}^2} = \frac{(l^2 - 3l + 3) s_{2,0}^2 + (l - l^2) s_{4,0} - 2ls_{2,0}s_{1,0}^2 + (4l - 4) s_{3,0}s_{1,0} + s_{1,0}^4}{(l-3)(l-2)(l-1)l} \quad (A6)$$

$$\widehat{\mu_{0,2}^2} = \frac{(l^2 - 3l + 3) s_{0,2}^2 + (l - l^2) s_{0,4} - 2ls_{0,2}s_{0,1}^2 + (4l - 4) s_{0,3}s_{0,1} + s_{0,1}^4}{(l-3)(l-2)(l-1)l} \quad (A7)$$

$$\begin{aligned} \widehat{\mu_{2,0}\mu_{0,2}} = & \frac{1}{(l-3)(l-2)(l-1)l} \left((l^2 - 3l + 1) s_{0,2}s_{2,0} + (l - l^2) s_{2,2} + \right. \\ & (2-l)s_{2,0}s_{0,1}^2 + (2l-2)s_{2,1}s_{0,1} + (2-l)s_{0,2}s_{1,0}^2 + \\ & \left. (2l-2)s_{1,0}s_{1,2} + s_{1,0}^2s_{0,1}^2 - 4s_{1,0}s_{1,1}s_{0,1} + 2s_{1,1}^2 \right) \end{aligned} \quad (A8)$$